# Spring Loaded Double Pantograph : Static Analysis

The static analysis provides a system of equations that can be solved to determine the internal forces in the SLDP mechanism for any given configuration and external load. Figure 1 illustrates the external forces acting at joints I and J. Due to the double-pantograph properties, links 7, 4, and 3 are in parallel, as are links 2, 8, and 1.

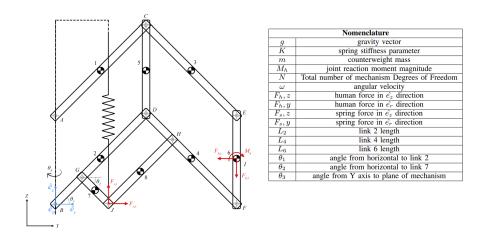


Figure 1: Definition of coordinate systems and kinematic parameters with accompanying nomenclature

## 1 Analysis of Link 2

We begin our analysis with link 2, as shown in Figure 2a. Applying force and moment balance yields:

$$F_{D2,y} + F_{G2,y} - F_{B2,y} = 0 (1)$$

$$F_{D2,z} + F_{G2,z} - F_{B2,z} = 0 (2)$$

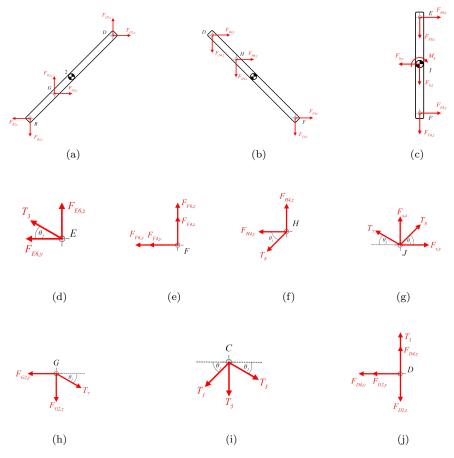


Figure 2: Forces on various links and joints

$$L_{GD} \sin \theta_1 F_{G2,y} - L_{GD} \cos \theta_1 F_{G2,z} - (L_{BG} + L_{GD}) \sin \theta_1 F_{B2,y} + (L_{BG} + L_{GD}) \cos \theta_1 F_{B2,z} = 0$$
(3)

where  $F_{xy,z}$  represents the force at point x in the y or z direction, and  $L_{xy}$  is the length between points x and y.

## 2 Analysis of Link 4

For link 4, as illustrated in Figure 2b, we have:

$$F_{F4,z} + F_{H4,z} + F_{D4,z} = 0 (4)$$

$$F_{F4,y} + F_{H4,y} + F_{D4,y} = 0 (5)$$

$$-L_{HF} \sin \theta_2 F_{H4,y} + L_{HF} \cos \theta_2 F_{H4,z} -(L_{DH} + L_{HF}) \sin \theta_2 F_{D4,y} +(L_{DH} + L_{HF}) \cos \theta_2 F_{D4,z} = 0$$
(6)

## 3 Analysis of Link 6

For link 6, which connects to the end effector, as illustrated in Figure 2c, we have:

$$F_{F6,z} + F_{H,z} + F_{E6,z} = 0 (7)$$

$$F_{F6,y} - F_{H,y} + F_{E6,y} = 0 (8)$$

$$F_{F6,y} - F_{E6,y} = \frac{2}{L_{EF}} M_H \tag{9}$$

where  $M_H$  is the moment applied at point H.

## 4 Analysis of Joints

Links 1, 3, 5, 7, and 8 can be modeled as two-force members with a constant axial force (denoted  $T_x$ ). For joint E:

$$F_{E6,z} + T_3 \sin \theta_2 = 0 \tag{10}$$

$$F_{E6,y} + T_3 \cos \theta_2 = 0 \tag{11}$$

For joint F:

$$F_{F4,z} + F_{F6,z} = 0 (12)$$

$$F_{F4,y} + F_{F6,y} = 0 (13)$$

For joint H:

$$F_{H4,z} - T_8 \sin \theta_1 = 0 \tag{14}$$

$$F_{H4,y} + T_8 \cos \theta_1 = 0 \tag{15}$$

For joint J, where the spring force  $F_s$  is applied:

$$F_{s,z} + T_7 \sin \theta_2 + T_8 \sin \theta_1 = 0 \tag{16}$$

$$-T_7 \cos \theta_2 + T_8 \cos \theta_1 + F_{s,y} = 0 \tag{17}$$

For joint G:

$$F_{G2,z} + T_7 \sin \theta_2 = 0 \tag{18}$$

$$-F_{G2,y} + T_7 \cos \theta_2 = 0 \tag{19}$$

For joint C:

$$T_5 + T_1 \sin \theta_1 + T_3 \sin \theta_2 = 0 \tag{20}$$

$$-T_1\cos\theta_1 + T_3\cos\theta_2 = 0\tag{21}$$

For joint D:

$$T_5 + F_{D4,z} - F_{D2,z} = 0 (22)$$

$$F_{D4,y} + F_{D2,y} = 0 (23)$$

The analysis presented here assumes static equilibrium. For dynamic situations, these equations would need to be augmented with inertial terms to account for accelerations of the links and the payload. Furthermore, the spring force  $F_s$  would be a function of the mechanism's configuration, typically modeled as a linear spring force:

$$F_s = k(\Delta l) \tag{24}$$

where k is the spring constant and  $\Delta l$  is the change in spring length from its rest position.